

Научно-координационная Сессия  
"Исследования неидеальной плазмы"

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УСТОЙЧИВОСТЬ КРИСТАЛЛА LiF  
В РАЗОГРЕТОМ ПЛОТНОМ СОСТОЯНИИ

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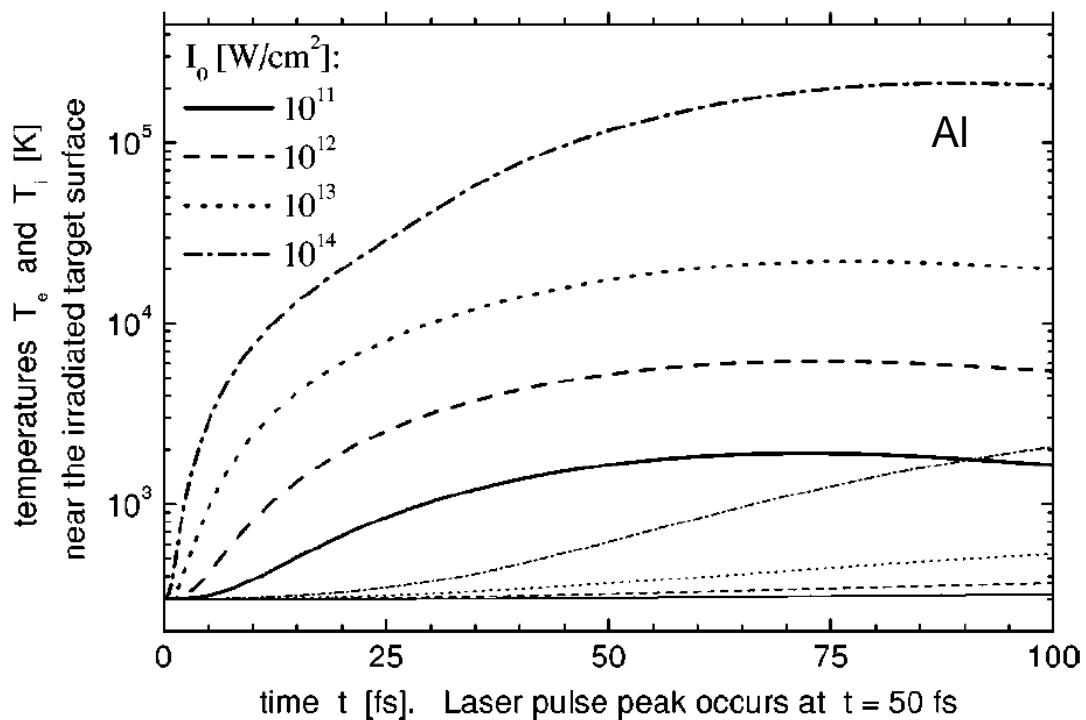
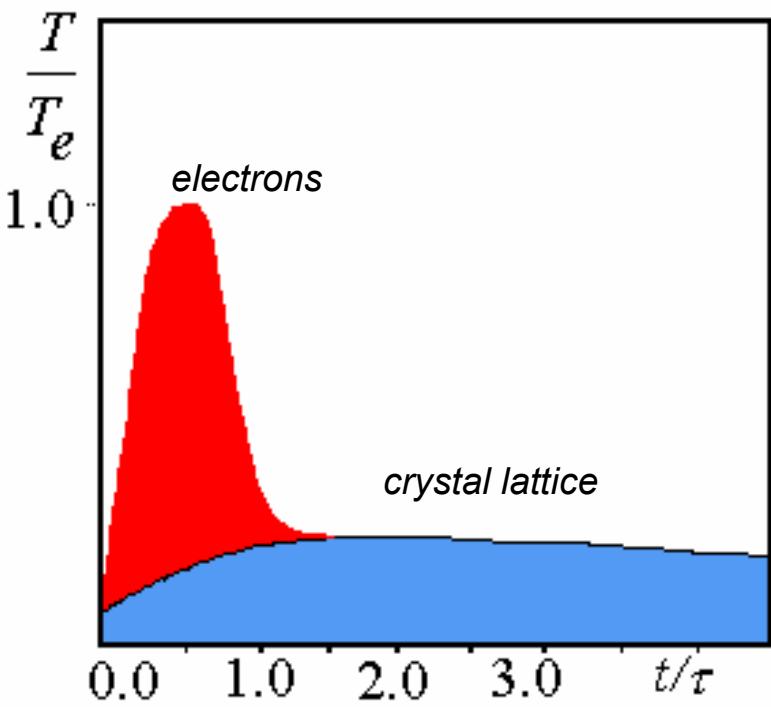
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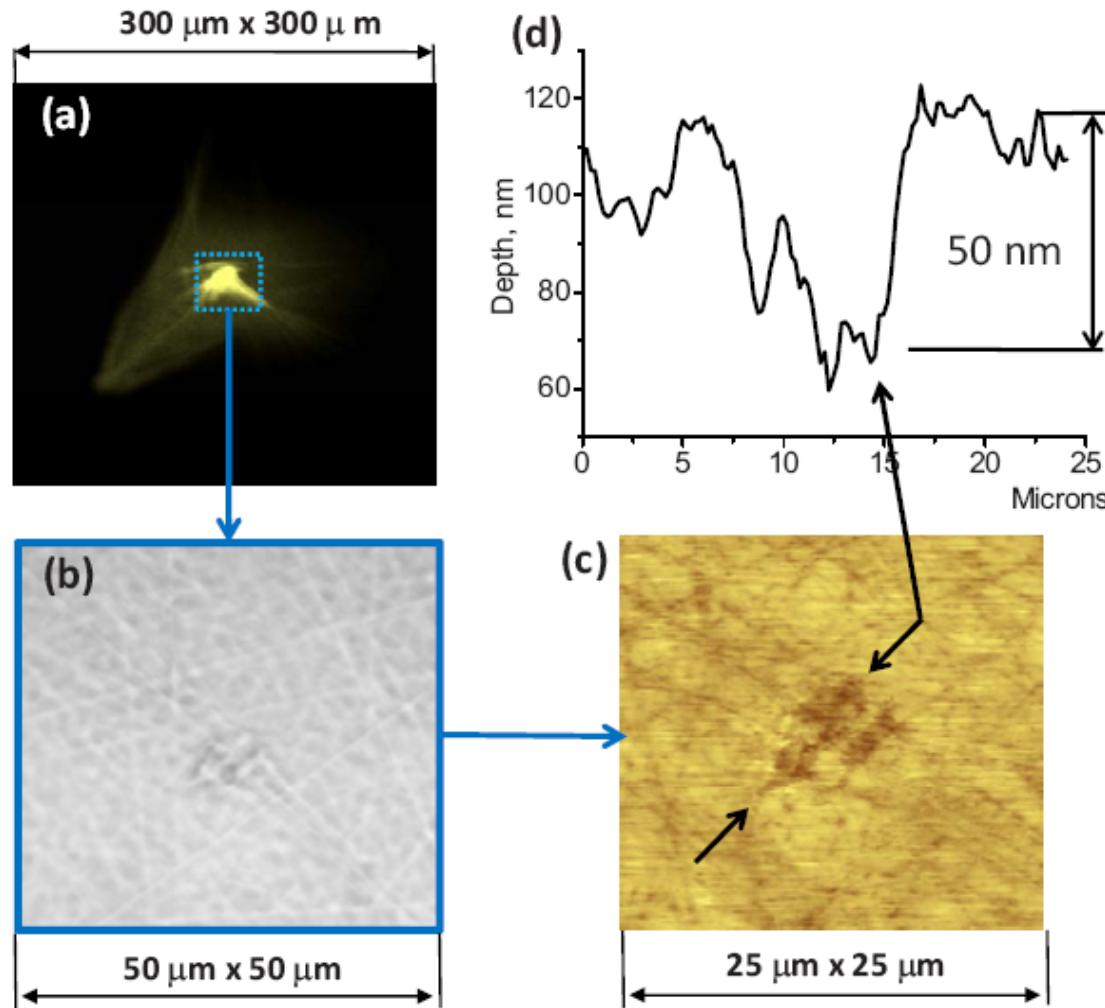
# Influence of the electronic system excitation on crystal lattice stability

# Dynamics of the crystal lattice after femtosecond pulse irradiation

Two-temperature system  
«cold» lattice – hot «electrons»



# Low-threshold ablation of dielectrics irradiated by picosecond soft x-ray laser pulses

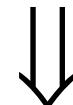


# Electronic density functional theory

# Density Functional Theory

N-electron  
wave function

$$\Psi(\vec{r}^N)$$



Hohenberg-Kohn theorem

One-electron  
density

$$n(\vec{r})$$



Kohn-Sham scheme

N one-electron  
orbitals

$$\Phi_1(\vec{r}), \dots, \Phi_N(\vec{r})$$

Connection:

$$N \int d\vec{r}_2 \dots d\vec{r}_N |\Psi(\vec{r}, \vec{r}_2 \dots \vec{r}_N)| = n(\vec{r}) = \sum_{i=1}^N |\Phi_i(\vec{r})|^2$$

# Density Functional Theory

$$E[n] = T_s[n] + V_H[n] + E_{xc}[n] + \int d\vec{r} n(\vec{r})v(\vec{r})$$

$$T_s[n] = -\frac{1}{2} \sum_{i=1}^N \langle \Phi_i | \nabla^2 | \Phi_i \rangle$$

$$\left( -\frac{1}{2} \nabla^2 + \frac{\delta V_H[n]}{\delta n(\vec{r})} + \underbrace{\frac{\delta E_{xc}[n]}{\delta n(\vec{r})}}_{+ v(\vec{r})} + v(\vec{r}) \right) \Phi_i = \epsilon_i \Phi_i$$

approximations to the exchange-correlation term  
LDA, GGA etc.

## T>0: Mermin functional

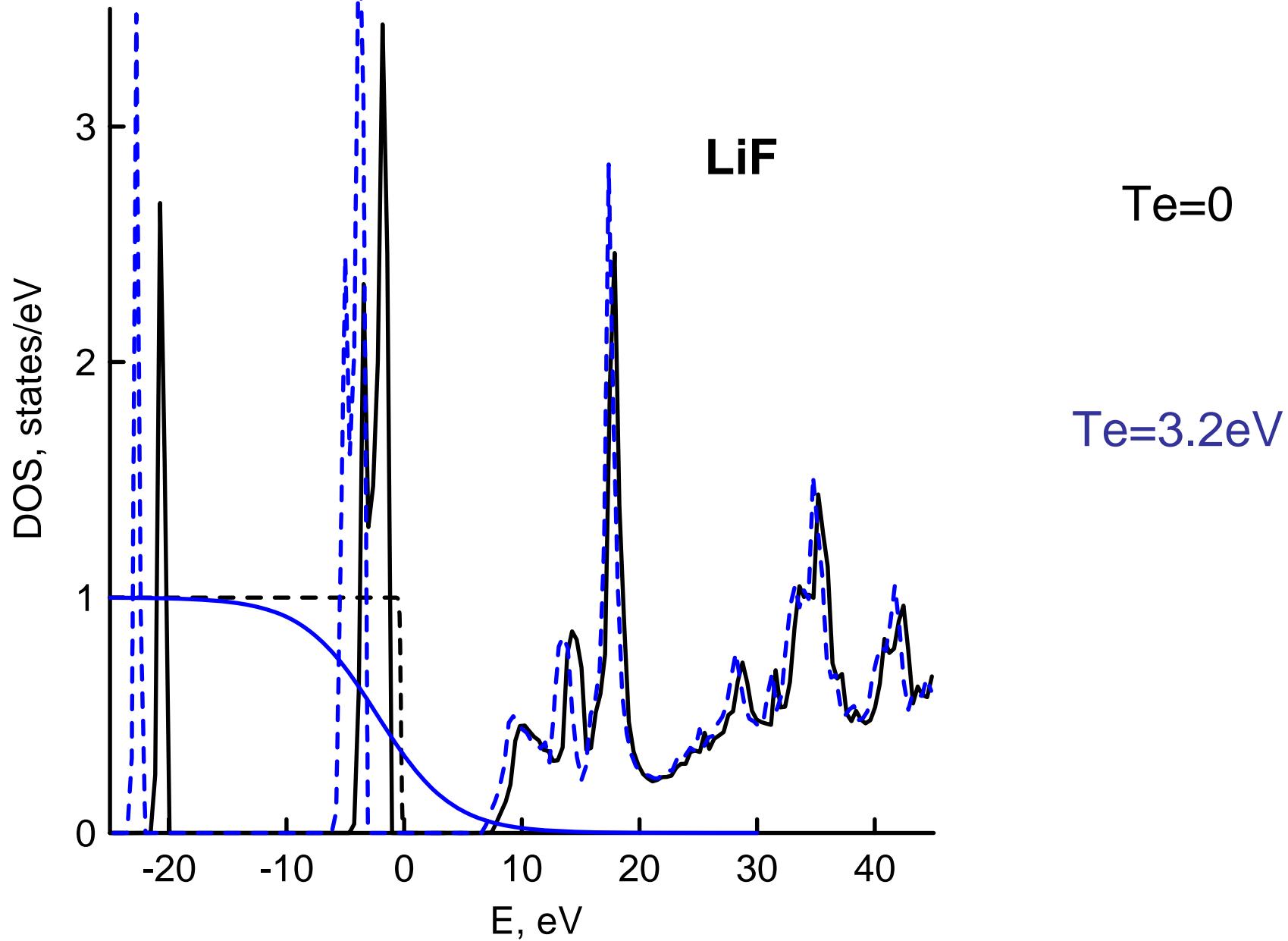
$$f_i(\beta) = \frac{1}{1 + \exp[\beta(\varepsilon_i - \mu)]}$$

$$n(\vec{r}) = \sum_{i=1}^N f_i(\beta) |\Phi_i(\vec{r})|^2$$

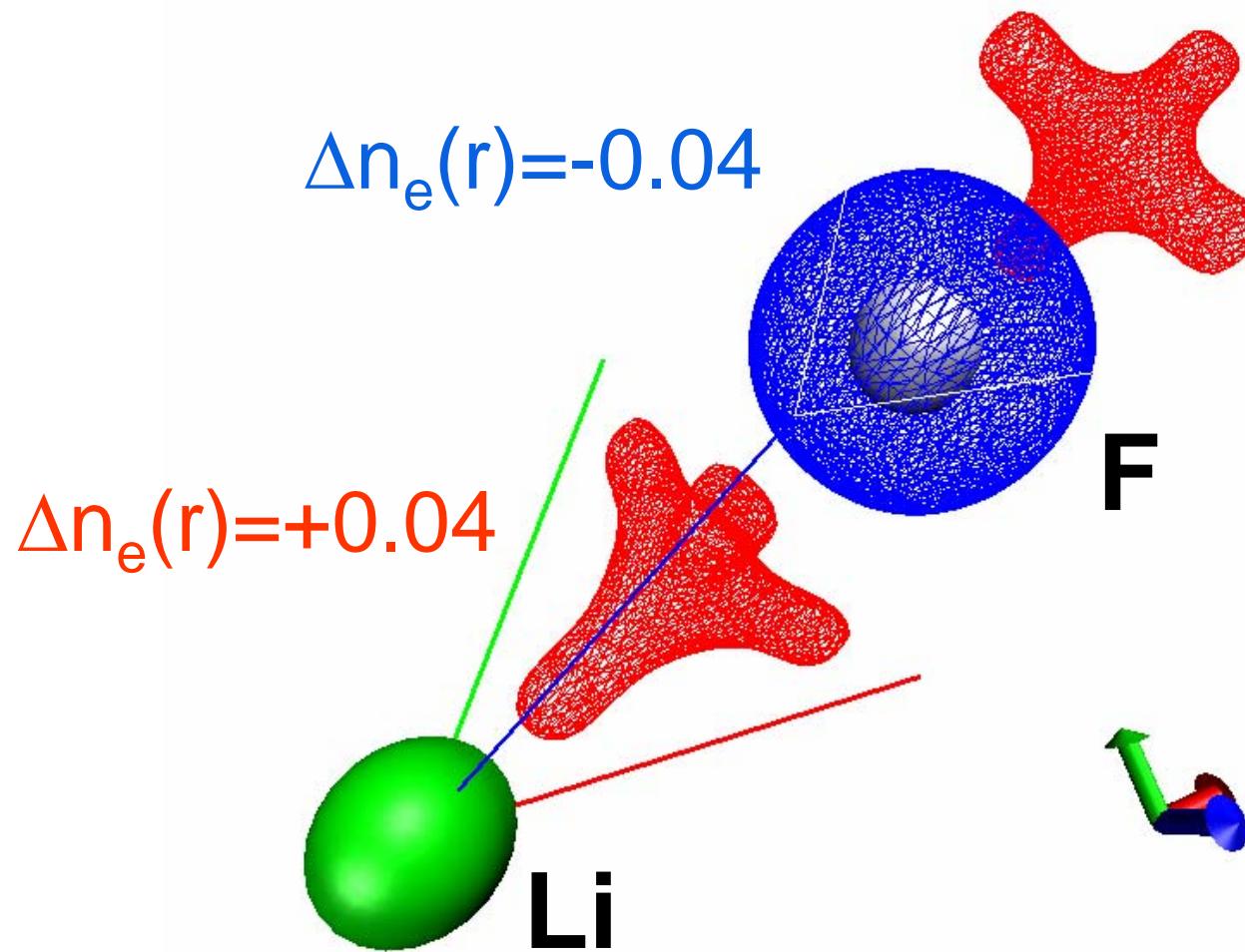
$$\Omega[n] = T_s[n] + V_H[n] + \Omega_{xc}[n] + \int d\vec{r} n(\vec{r}) v(\vec{r})$$

Mermin, 1965

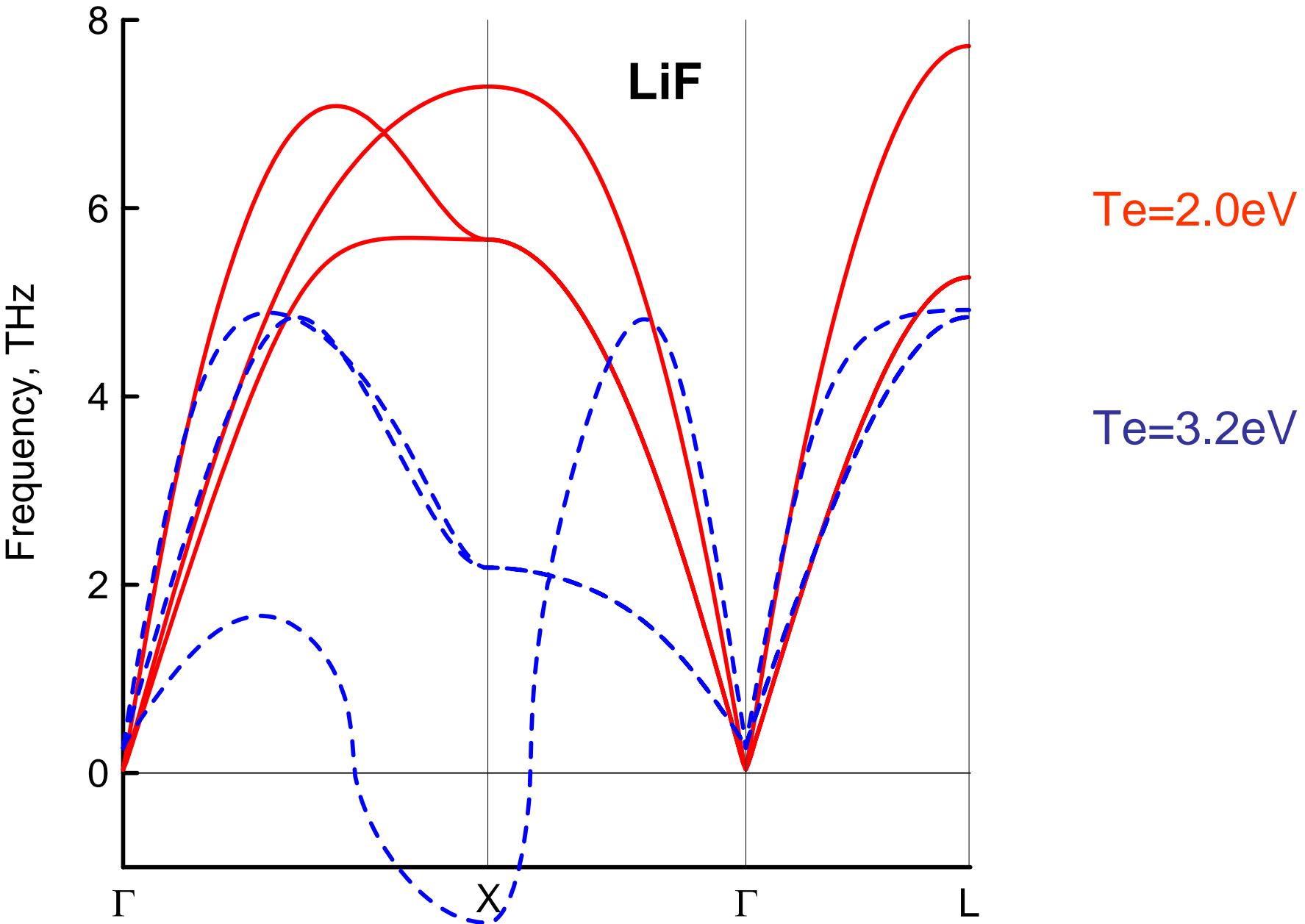
# LiF: electron DOS



# Redistribution of the electron density due to the heating of the electron subsystem in the LiF crystal unit cell



# Phonon dispersion in LiF



# Conclusions

**The excitation of the electronic subsystem results in the loss of mechanical stability of the fcc LiF lattice that is manifested as an appearance of the soft acoustic phonon mode. The corresponding redistribution of the electronic density implies that the originally strongly ionic interatomic interaction becomes more of covalent character with the rise of electronic temperature.**